



UDC 546.62'31+546.714

## LUMINESCENT PROPERTIES OF Mn<sup>4+</sup>-DOPED α-Al<sub>2</sub>O<sub>3</sub> OBTAINED BY COMBUSTION METHOD

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Received 10 June 2022; accepted 29 November 2022; available online 26 January 2023

### Abstract

The Mn<sup>4+</sup>-doped α-Al<sub>2</sub>O<sub>3</sub> with particle size of 70–600 nm was obtained by combustion method. Based on the results of luminescence measurements, the crystal-field strength (*Dq*) and Racah parameters (*B*, *C*) for Mn<sup>4+</sup> in α-Al<sub>2</sub>O<sub>3</sub> were determined using a pure electronic transition approach. The obtained values of *Dq* (1898 cm<sup>-1</sup>) and nephelauxetic parameter β<sub>1</sub> (0.95) for Mn<sup>4+</sup> ions in α-Al<sub>2</sub>O<sub>3</sub> are consistent with those of Mn<sup>4+</sup> in other oxide compounds. In particular, the *Dq/B* value of 2.77 indicates a strong crystal field environment of the Mn<sup>4+</sup> ions in the α-Al<sub>2</sub>O<sub>3</sub> lattice. A comparison with literature data for Cr<sup>3+</sup> in α-Al<sub>2</sub>O<sub>3</sub> was also carried out. The composites of general formula α-Al<sub>2</sub>O<sub>3</sub> : Mn<sup>4+</sup>, Mg<sup>2+</sup>/Gd<sub>3</sub>Al<sub>5</sub>O<sub>12</sub>:Ce<sup>3+</sup> had been also obtained. It is shown that these materials demonstrate the intense broadband emission with maxima at about 585 and 678 nm. The broad band with a maximum at 585 nm is caused by the 5d→4f transitions of Ce<sup>3+</sup> ions in Gd<sub>3</sub>Al<sub>5</sub>O<sub>12</sub>, whereas the band at 678 nm is caused by the Mn<sup>4+</sup> <sup>2</sup>E<sub>g</sub>→<sup>4</sup>A<sub>2g</sub> transitions in α-Al<sub>2</sub>O<sub>3</sub>. The emission color changes from yellow to deep red with increasing content of α-Al<sub>2</sub>O<sub>3</sub>:Mn<sup>4+</sup>, Mg<sup>2+</sup>, and the luminescence quantum efficiency of the composites was found as high as 0.60.

**Keywords:** aluminum oxide; combustion synthesis; Mn<sup>4+</sup>; luminescence; Racah parameters; luminescence composites.

## ЛЮМІНЕСЦЕНТНІ ВЛАСТИВОСТІ α-Al<sub>2</sub>O<sub>3</sub> : Mn<sup>4+</sup>, ОДЕРЖАНОГО МЕТОДОМ ГОРІННЯ

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### Анотація

Методом горіння одержано легований Mn<sup>4+</sup> α-Al<sub>2</sub>O<sub>3</sub> з розміром частинок 70–600 нм. На підставі результатів люмінесцентних вимірювань в рамках наближення чисто електронних переходів розраховано силу кристалічного поля (*Dq*) та параметри Рака (*B*, *C*) для Mn<sup>4+</sup> в α-Al<sub>2</sub>O<sub>3</sub>. Одержані значення *Dq* (1898 см<sup>-1</sup>) та нефелауксетичного параметра β<sub>1</sub> (0.95) для іонів Mn<sup>4+</sup> в α-Al<sub>2</sub>O<sub>3</sub> узгоджуються зі значеннями цих параметрів для Mn<sup>4+</sup> в інших оксидних сполуках. Зокрема, величина *Dq/B* = 2.77 вказує, що іони Mn<sup>4+</sup> займають в ґратці α-Al<sub>2</sub>O<sub>3</sub> октаедричні позиції Al з великою силою кристалічного поля. Проведено порівняння з літературними даними для іонів Cr<sup>3+</sup> в α-Al<sub>2</sub>O<sub>3</sub>. Одержано композити загальної формули α-Al<sub>2</sub>O<sub>3</sub>:Mn<sup>4+</sup>, Mg<sup>2+</sup>/Gd<sub>3</sub>Al<sub>5</sub>O<sub>12</sub>:Ce<sup>3+</sup>. Показано, що ці матеріали демонструють інтенсивне широкомуглове випромінювання з максимумами при ~585 і 678 нм. Широка смуга з максимумом при 585 нм зумовлена переходами 5d→4f іонів Ce<sup>3+</sup> в Gd<sub>3</sub>Al<sub>5</sub>O<sub>12</sub>, тоді як смуга при 678 нм зумовлена переходами <sup>2</sup>E<sub>g</sub>→<sup>4</sup>A<sub>2g</sub> Mn<sup>4+</sup> в α-Al<sub>2</sub>O<sub>3</sub>. Колір випромінювання змінюється від жовтого до насичено-червоного із збільшенням вмісту α-Al<sub>2</sub>O<sub>3</sub>:Mn<sup>4+</sup>, Mg<sup>2+</sup>, а квантова ефективність люмінесценції композитів досягає 0.60.

**Ключові слова:** оксид алюмінію; синтез методом горіння; Mn<sup>4+</sup>; люмінесценція; параметри Рака; люмінесцентні композити.

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doi: 10.15421/jchemtech.v30i4.259866

## Introduction

In recent years, crystals doped with tetravalent manganese ions ( $Mn^{4+}$ ) have been adopted in various fields, such as general lighting [1; 2], plant cultivation light-emitting diodes (LED) [3] and LED displays [4]. The luminescence properties of  $Mn^{4+}$  in inorganic materials have been examined in a number of papers [2; 5–7]. It is known that due to a high effective positive charge,  $Mn^{4+}$  tends to occupy octahedral sites with large crystal-field strength, and only transitions from the  ${}^2E_g$  state to the  ${}^4A_{2g}$  ground state are observed in the emission spectra. The  ${}^2E_g \rightarrow {}^4A_{2g}$  electronic transitions are spin-forbidden, but they become partially allowed due to mixing of the  ${}^2E_g$  and  ${}^4T_{2g}$  states by spin-orbit interaction. Typically, the emission spectra of  $Mn^{4+}$  consist of narrow zero-phonon lines (ZPL or R-lines) and corresponding vibronic sidebands. It is shown that the energy position of the  ${}^2E_g$  state is nearly independent on the crystal-field strength, but it is substantially dependent on the covalency of the  $Mn^{4+}$ -ligand bond. The influence of the main lattice on the luminescence properties of transition metal ions with the  $3d^3$ -configuration ( $Cr^{3+}$ ,  $Mn^{4+}$ ) is usually described by a number of parameters, such as the crystal-field strength ( $Dq$ ) and the Racah parameters ( $B$ ,  $C$ ) in conjunction with the Tanabe-Sugano diagram [2; 6; 7]. The crystal-field strength describes an interaction of the d-electrons with the electrostatic field of the ligands, while the Racah parameters characterize the effects of inter-electronic repulsion.

Because of the covalent interaction between  $Mn^{4+}$  and anions (or ligands), the Racah parameters are reduced in comparison with the free ion values ( $B_0 = 1160 \text{ cm}^{-1}$ ,  $C_0 = 4303 \text{ cm}^{-1}$ ), and this effect can be expressed by the so called nephelauxetic parameter  $\beta_1$  [5; 6]:

$$\beta_1 = \sqrt{\left(\frac{B}{B_0}\right)^2 + \left(\frac{C}{C_0}\right)^2} \quad (1)$$

In 2015 it has been reported that energy of the  $Mn^{4+}$   ${}^2E_g$  state ( $E({}^2E_g)$ ) is a linear function of the parameter  $\beta_1$  for different kinds of host materials, including fluorides, silicates and also numerous aluminates [5; 6]. In this regard, it seems rather surprising that one of few exceptions from this trend is thermodynamically stable polymorph ( $\alpha$ ) of  $Al_2O_3$  [5; 6]. The reasons for this have not been clarified. It should be noted that in  $\alpha$ - $Al_2O_3$ , each  $Al^{3+}$  is surrounded by a trigonally distorted octahedron of six oxygen atoms.

Very recently the authors [8] have reported the optical properties of red emitting ceramics of

composition  $\alpha$ - $Al_2O_3:Mn^{4+}$ ,  $Mg^{2+}$ . Based on the results obtained, the calculated ratio of  $Dq$  to  $B$  for the  $Mn^{4+}$  appeared as 1.74. This value of  $Dq/B$  implies a weak crystal field surrounding of the  $Mn^{4+}$  ions in  $\alpha$ - $Al_2O_3$  host [8]. This conclusion seems to be doubtful, since, as it was mentioned above, the  $Mn^{4+}$  ions are always experiencing the strong crystal field, i.e. the  $Dq/B$  value should be  $\geq 2.2$  [2; 6; 7].

It should be also noted that  $\alpha$ - $Al_2O_3$  doped with  $Mn^{4+}$  appears as a promising phosphor for non-contact luminescence thermometry [9]. In view of the importance of  $Mn^{4+}$ -doped  $\alpha$ - $Al_2O_3$  as a functional material and also as a model system to analyse different aspects of spectroscopy of the  $3d^3$  ions ( $Mn^{4+}$ ,  $Cr^{3+}$ ,  $V^{2+}$ ) in inorganic compounds, the aim of this study was finding out the reasons for deviation of the data for  $Mn^{4+}$ -doped  $\alpha$ - $Al_2O_3$  from linear dependence of energy of the  ${}^2E_g$  state on  $\beta_1$  [5; 6]. Besides, the present work was motivated by an interest to obtain luminescent composites based on  $Al_2O_3:Mn^{4+}$  and, in particular, ones of general formula  $Al_2O_3:Mn^{4+}/Gd_3Al_5O_{12}:Ce^{3+}$ . Gadolinium aluminum garnet  $Gd_3Al_5O_{12}$  (GdAG) doped with  $Ce^{3+}$  ions is known as an effective phosphor for medium power white LEDs [10; 11]. Also, the development of nanocomposites based on GdAG: $Ce^{3+}$  has been recently reported for imaging and photodynamic therapy of cancer cells [12].

## Experimental

A powder sample of Mn-doped  $Al_2O_3$  with particles sizes of 10–70 nm was prepared by a gas-dispersed synthesis (GDS). This method is based on the combustion of powdered metals due to exothermic oxidizing reactions between them and a gaseous oxidizer (typically  $O_2$ ) [13], and it is shown to provide favorable morphology of  $Al_2O_3$  particles and good dispersibility of the final product in various media [13–15]. The details of experimental apparatus used for the GDS of nanosized Mn-doped  $Al_2O_3$  were described previously [13]. As it was also found earlier [16], the resulting powders consist of a mixture of transition aluminas ( $\delta^*$ ,  $\delta$  and  $\theta$ -polymorphs), and a part of Mn ions exists in the oxidation state +2 and occupies tetrahedral positions in  $\delta^*$ - $Al_2O_3$ , causing a broadband luminescence with a maximum at  $\sim 520 \text{ nm}$ . It is shown that additional annealing in air at temperatures  $\geq 1130^\circ\text{C}$  results in the formation of stable  $\alpha$ -polymorph. The  $\delta^*, \theta \rightarrow \alpha$ - $Al_2O_3$  phase transition is followed by oxidation of  $Mn^{2+}/Mn^{3+}$  ions and the formation of some amount of manganese ions in the oxidation

state +4 on octahedral Al positions [16]. In  $\text{Al}_2\text{O}_3$ , the stabilization of  $\text{Mn}^{4+}$  can be provided through the formation of cation (Al) vacancies [17; 18] or by the introduction into the host lattice of charge-compensating ions, such as  $\text{Mg}^{2+}$  [8; 19].

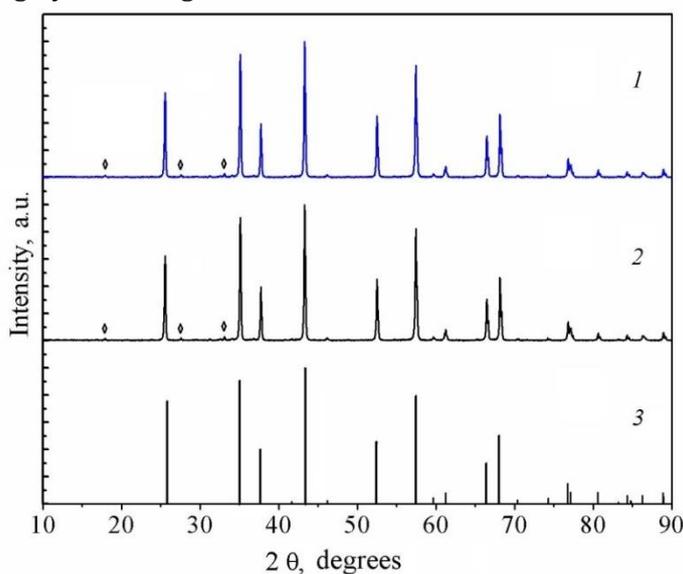
To obtain Mn-doped  $\alpha\text{-Al}_2\text{O}_3$  with a nominal Mn concentration of 0.1 at.%, the appropriate amount of  $\text{MnCl}_2$  (99.99 %) was dissolved in ethanol. Then, the solution was gradually added, while thoroughly mixing, to calculated amounts of Al powder (purity 99.7 %) with an average size of 4.8  $\mu\text{m}$ . After drying in air, the as-obtained mixture was dispersed in nitrogen and injected through an inner tube into oxygen stream. After ignition by an external source a stable two-phase diffusion dust flame was observed. The condensed products of aluminum combustion had been collected using a polyethylene terephthalate filter. Some part of the sample obtained was then used to prepare double doped  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}$  and luminescence composites  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}/\text{GdAG}:\text{Ce}^{3+}$ . For this purpose the calculated amount of magnesium chloride ( $\text{MgCl}_2 \cdot 6\text{H}_2\text{O}$ ) was dissolved in ethanol and added to the nanosized Mn-doped  $\text{Al}_2\text{O}_3$  to improve the luminescence efficiency of final products. This reagent was used in the amount corresponding to the atomic Mg/Mn ratio of 1. After drying at 110 °C all the samples were annealed in air at the temperature of 1300 °C during 6 hours.

Luminescence composites  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}/\text{GdAG}:\text{Ce}^{3+}$ , with different weight ratios of the components, had been obtained by thoroughly mixing of  $\alpha\text{-}$

$\text{Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}$  and  $\text{GdAG}:\text{Ce}^{3+}$  powders in ethanol, followed by drying and annealing in air at the temperature of 1300 °C during for 2 hours.  $\text{Ce}^{3+}$ -doped gadolinium aluminum garnet  $\text{GdAG}:\text{Ce}^{3+}$  with a nominal  $\text{Ce}^{3+}$  concentration of 1 at.% was previously synthesized by the coprecipitation of rare earth metals and aluminum hydroxides and their further thermal decomposition [11]. X-ray diffraction (XRD) studies had been carried out on a Rigaku Ultima IV diffractometer using Cu K $\alpha$  radiation ( $\lambda = 1.5418 \text{ \AA}$ ). Transmission electron microscopy (TEM) images of the samples had been obtained on a Philips EM-400 transmission electron microscope. The emission and excitation spectra in UV-visible region were recorded at 293 and 77 K using a Fluorolog FL-3-22 (Horiba Jobin Yvon) spectrofluorometer equipped with a xenon lamp. Diffuse reflection spectra had been obtained on a Perkin-Elmer Lambda 9 spectrometer.

## Results and discussion

Fig. 1 shows the XRD patterns of Mn-doped  $\alpha\text{-Al}_2\text{O}_3$  and  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}$ . One can see that the XRD patterns of the samples annealed at 1300 °C for 6 h match well with data from JCPDS database for pure  $\alpha\text{-Al}_2\text{O}_3$  (JCPDS Card 10-173), but trace amounts of  $\theta\text{-Al}_2\text{O}_3$  are also detected. It should be noted that the introduction of Mn and Mg is not accompanied by any changes in the XRD patterns. It is an expected result because of the relatively small concentrations of Mn and Mg in the samples. TEM image of as-prepared  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}$  is presented in Fig. 2.



**Fig. 1.** Comparison of X-ray diffraction patterns of as-prepared Mn-doped  $\alpha\text{-Al}_2\text{O}_3$  (1) and  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}$  (2) with data from JCPDS database (JCPDS Card 10-173) for pure  $\alpha\text{-Al}_2\text{O}_3$  (3). Reflections associated with the presence of trace amounts ( $\leq 2\%$ ) of  $\theta\text{-Al}_2\text{O}_3$  are shown by symbols (o)

As it can be seen, the sample consists of submicron-size (200–600 nm) aggregates of vermicular shape, although some amount of nanosized particles is also present as individual crystallites. It should be noted that electron microscopy data obtained in the present study did not also reveal any significant effect of Mn (Mg) on the particle size distribution of the final products. The Kubelka-Munk transformed diffuse reflection spectrum of as-prepared  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>:Mn appeared to be similar to those reported for Mn<sup>3+</sup>-doped  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> [20]. It consists of wide overlapping bands in the 400–620 nm area. There are two local maxima at ~490 and 537 nm. Analogous absorption spectra had been also obtained for Mn<sup>3+</sup>-doped

yttrium aluminates Y<sub>3</sub>Al<sub>5</sub>O<sub>12</sub> [21] and YAlO<sub>3</sub> [22], and they had been interpreted as a result of <sup>5</sup>E<sub>g</sub>→<sup>5</sup>T<sub>2</sub> transitions in Mn<sup>3+</sup> ions at the aluminum octahedral (Al<sub>oct</sub>) positions. The substantial half-width (~4000 cm<sup>-1</sup>) of Mn<sup>3+</sup> absorption band in Y<sub>3</sub>Al<sub>5</sub>O<sub>12</sub> was attributed to Jahn-Teller splitting of both <sup>5</sup>E<sub>g</sub>, <sup>5</sup>T<sub>2</sub> states of Mn<sup>3+</sup> [21]. Thus, the broad absorption bands in the diffuse reflection spectrum of as-prepared  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>:Mn are mainly caused by the <sup>5</sup>E<sub>g</sub>→<sup>5</sup>T<sub>2</sub> transitions of Mn<sup>3+</sup> ions occupying the Al<sub>oct</sub> positions.

The emission and luminescence excitation spectra of the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>:Mn at 293 and 77 K are compared in Figs. 3 and 4.

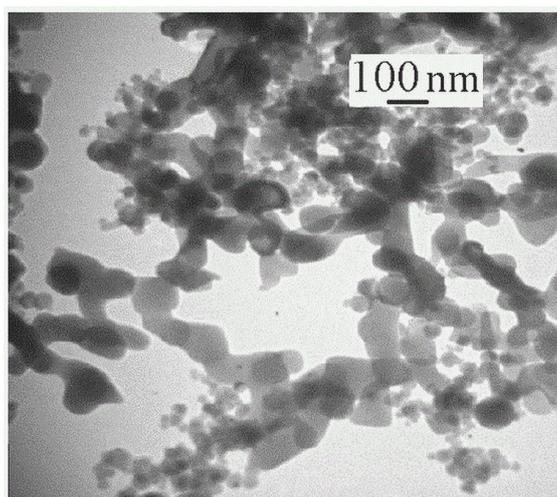


Fig. 2. TEM image of as-prepared  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>:Mn

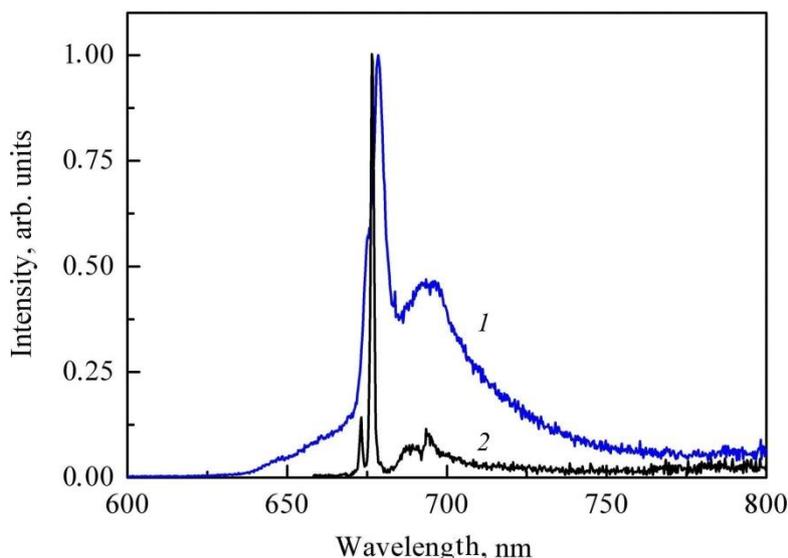


Fig. 3. Comparison of the normalized emission spectra of Mn<sup>4+</sup> in as-prepared  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>:Mn<sup>3+/4+</sup> upon excitation with  $\lambda_{exc} = 460$  nm at 293 K (1) and 77 K (2)

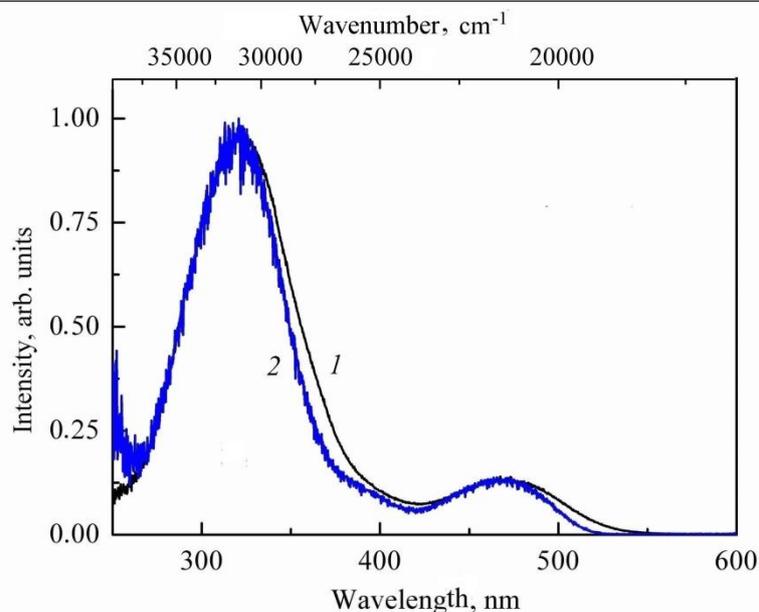


Fig. 4. Comparison of the excitation spectra for the  $\text{Mn}^{4+}$  emission ( $\lambda_{\text{em}} = 678 \text{ nm}$ ) in as-prepared  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{3+/4+}$  recorded at 293 K (1) and 77 K (2)

At the room temperature the emission spectrum contains an intense band at 678 nm with a long-wavelength tail extending up to 770 nm. There is no doubt about its nature, because the identical spectrum was reported earlier for  $\text{Mn}^{4+}$ -doped  $\alpha\text{-Al}_2\text{O}_3$  [19; 20], and it was interpreted as a result of the  ${}^2\text{E}_g \rightarrow {}^4\text{A}_{2g}$  transition and vibronic transitions, respectively. A deviation from perfect octahedral symmetry and spin-orbit interaction are known to split the  ${}^2\text{E}_g$  state into two components, and they result in so-called  $\text{R}_{1,2}$ -lines luminescence, which is caused with the parity- and spin-forbidden  ${}^2\text{E}_g \rightarrow {}^4\text{A}_{2g}$  transitions of  $\text{Mn}^{4+}$  ions. It is seen from Fig. 3 that at 77 K the emission spectrum consists of two relatively narrow lines at 676.4 nm ( $14784 \text{ cm}^{-1}$ ) and 672.9 nm ( $14861 \text{ cm}^{-1}$ ), known as  $\text{R}_1$  and  $\text{R}_2$ -lines, and vibronic sideband at 683–710 nm.

It should be noted that in the 77–293 K range the sample did not show any emission that could be attributed to either  $\text{Mn}^{3+}$  or  $\text{Mn}^{2+}$  ions. This agrees with the conclusion of the authors [20] about the absence of any luminescence from  $\text{Mn}^{3+}$  ions in  $\alpha\text{-Al}_2\text{O}_3$  even at 4.2 K, which was explained by a strong Jahn-Teller effect.

It can be seen from Fig. 4 that at 293 K the excitation spectrum for the luminescence of  $\text{Mn}^{4+}$  ions in  $\alpha\text{-Al}_2\text{O}_3$  consists of a broad intense band with a maximum at 323 nm ( $30960 \text{ cm}^{-1}$ ) and a less intense one at 476 nm ( $21008 \text{ cm}^{-1}$ ). The latter is obviously connected with the  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{2g}$  transition of  $\text{Mn}^{4+}$  ions, while the assignment of the former band is not unambiguous. In Ref. [8] the broad band at about 316 nm was attributed to the  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{1g}$  transition of  $\text{Mn}^{4+}$  ions, whereas the

authors [20] assigned this feature to a charge-transfer transition from 2p orbitals of  $\text{O}^{2-}$  to the  $\text{Mn}^{4+}$  vacant orbitals. An analogous band at 317 nm was shown to be present in the luminescence excitation spectra for  $\text{Mn}^{4+}$  ions in  $\text{GdAlO}_3$ . Based on the calculations for  $\text{Mn}^{4+}$  energy levels in  $\text{GdAlO}_3$  in limits of the exchange charge model of crystal-field theory, it was interpreted as a superposition of the  $\text{O}^{2-} \rightarrow \text{Mn}^{4+}$  charge transfer band and a band caused by the  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{1g}$  transition of  $\text{Mn}^{4+}$  ions located in  $\text{Al}_{\text{oct}}$  positions [23]. The excitation spectrum of the  $\text{Mn}^{4+}$  luminescence in as-prepared  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{3+/4+}$  recorded at 77 K is similar to that at the room temperature. But in addition to the bands at 323 and 476 nm, an arm in the 380–420 nm range is distinctly observed, suggesting a superposition character of the 323 nm band.

Usually, instead of zero-phonon line (ZPL) energies, the energies relating to the maxima of  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{1g}$ ,  ${}^4\text{T}_{2g}$  excitation bands of  $\text{Mn}^{4+}$  are used to determine the values of  $Dq$ ,  $B$  and  $C$ . However, this often leads to a big error or uncertainty in the estimation of the Racah parameters. In our opinion, this is also the main reason for the unreliable  $Dq/B$  value of 1.74 obtained in Ref. [8] for  $\text{Mn}^{4+}$  in  $\alpha\text{-Al}_2\text{O}_3$ , and the essential deviation of the data for  $\text{Mn}^{4+}$ -doped  $\alpha\text{-Al}_2\text{O}_3$  from the linear dependence of  $E({}^2\text{E}_g)$  on  $\beta_1$  [5; 6]. Meanwhile, even for the ordered crystals, an accurate determination of the ZPL energies of the  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{1g}$ ,  ${}^4\text{T}_{2g}$  transitions is very difficult because of the crystal-field splitting of  ${}^4\text{T}_{1g}$ ,  ${}^4\text{T}_{2g}$  levels, and an interaction of  $\text{Mn}^{4+}$  with lattice vibration modes. To overcome this problem a new approach based

on the ligand field theory had been proposed recently [7; 24]. In this approach the following

$$E(^4T_{2g}) = 10 Dq \quad (2);$$

$$B = 6.18Dq - \frac{1}{2}[(12.36Dq)^2 - 2.22E(^2E_g) Dq]^{1/2} \quad (3);$$

$$\frac{C}{B} = 4.7 \quad (4);$$

$$E(^4T_{1g})_{ZPL} = 7.5B + 15Dq - \frac{Dq}{2} \left[ \left( \frac{15B}{Dq} + 10 \right)^2 - \frac{480B}{Dq} \right]^{1/2} \quad (5);$$

Since in the case of  $\alpha\text{-Al}_2\text{O}_3\text{:Mn}^{4+}$  the fine structure of interest here turns out to be obscured, the wavelength where the excitation intensity has dropped to  $\sim 3\%$  on the long wavelength side of the  $^4A_{2g} \rightarrow ^4T_{2g}$  excitation band was used in the present study as a rough estimate for  $E(^4T_{2g})_{ZPL}$  of the  $\text{Mn}^{4+}$  ions in the  $\alpha\text{-Al}_2\text{O}_3$ . Obviously, the ZPL of the  $^4A_{2g} \rightarrow ^4T_{2g}$  excitation band should be located near its long-wavelength tail, and the used assumption conforms to the results of calculations

expressions are applied to calculate the Racah and crystal-field related parameters:

based on the Franck-Condon analysis model for numerous  $\text{Mn}^{4+}$ -doped inorganic compounds [2; 7; 24]. From the positions of the  $R_1$ -line ( $14784 \text{ cm}^{-1}$ ) and  $^4A_{2g} \rightarrow ^4T_{2g}$  excitation band in the spectra at 77 K (Figs. 3 and 4), the Racah and crystal-field related parameters for  $\text{Mn}^{4+}$  were calculated using Eqs. (2)–(5). The obtained data are presented in the Table and compared with those for  $\text{Cr}^{3+}$  ions in  $\alpha\text{-Al}_2\text{O}_3$  [25; 26].

Table

Comparison of the crystal-field strength  $Dq$ , Racah parameters and nephelauxetic parameter  $\beta_1$  for  $\text{Mn}^{4+}$  and  $\text{Cr}^{3+}$  ions in  $\alpha\text{-Al}_2\text{O}_3$  (T= 77 K)

Ion	$E(^2E_g)$ ( $\text{cm}^{-1}$ )	$Dq$ ( $\text{cm}^{-1}$ )	$B$ ( $\text{cm}^{-1}$ )	$C$ ( $\text{cm}^{-1}$ )	$Dq/B$	$\beta_1$	Ref.
$\text{Cr}^{3+}$	14424	1676	669	3144	2.50	1.09	[25,26]
$\text{Mn}^{4+}$	14784	1898	684	3215	2.77	0.95	This work

As it can be seen from the Table, the crystal-field strength for  $\text{Mn}^{4+}$  in the  $\alpha\text{-Al}_2\text{O}_3$  is  $1898 \text{ cm}^{-1}$ , which exceeds that for  $\text{Cr}^{3+}$  ( $Dq = 1676 \text{ cm}^{-1}$ ). This is the expected result because of the bigger formal charge and the smaller ionic radius of  $\text{Mn}^{4+}$  compared to  $\text{Cr}^{3+}$ . Since the electronegativity of  $\text{Mn}^{4+}$  ( $\chi = 1.912$ ) is higher than that of  $\text{Cr}^{3+}$  ( $\chi = 1.587$ ) [27], it can be expected that the covalency of the  $\text{Mn}^{4+}\text{-O}$  bond should be bigger than that of the  $\text{Cr}^{3+}\text{-O}$  bond. This is also conformed to the obtained data: the  $\beta_1$  value of 0.95 for  $\text{Mn}^{4+}$  in the  $\alpha\text{-Al}_2\text{O}_3$  is found to be smaller than that ( $\beta_1 = 1.09$ ) reported previously for  $\text{Cr}^{3+}$  ions in this compound [9; 25]. The reliability of the results obtained for

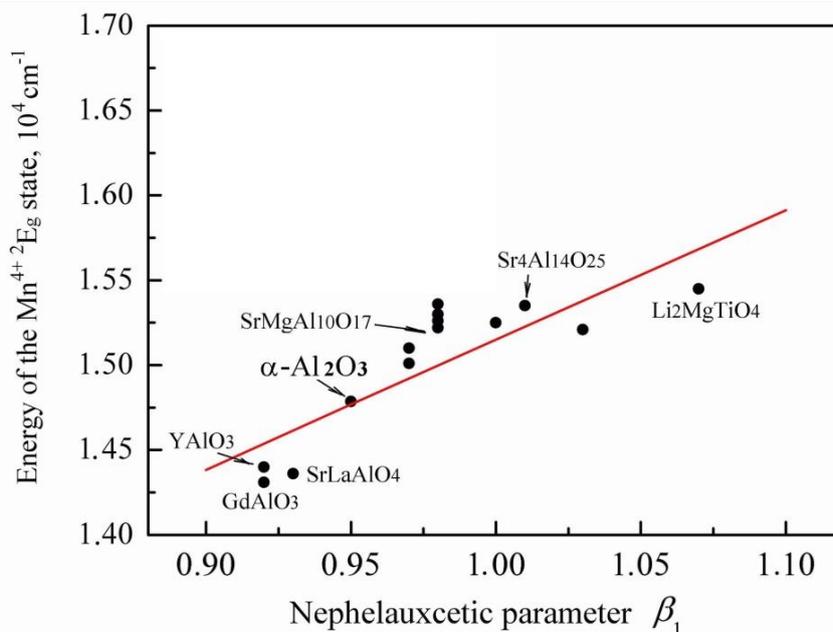
the  $\text{Mn}^{4+}$ -doped  $\alpha\text{-Al}_2\text{O}_3$  is confirmed by the fact that the calculated excited state energies divided by  $B$  and plotted against  $Dq/B$ , fall on the Tanabe-Sugano diagram well. In this context, it is worth noting that the  $Dq/B$  value of 2.77 (see the Table) indicates a strong crystal field surrounding of the  $\text{Mn}^{4+}$  ions in the  $\alpha\text{-Al}_2\text{O}_3$  lattice.

Fig. 5 shows the dependence of energy of the  $\text{Mn}^{4+} \ ^2E_g$  state on the nephelauxetic parameter  $\beta_1$  for a large group of oxide compounds. With the exception of  $\alpha\text{-Al}_2\text{O}_3$ , the  $E(^2E_g)$  and  $\beta_1$  values for these compounds were taken from Ref. [2] and the solid line represents the least-squares fit of the data by the expression proposed also in [2]:

$$E(^2E_g) = 7.63 \times 10^3 \beta_1 + 7.52 \times 10^3 \quad (\text{in } \text{cm}^{-1} \text{ units}) \quad (6).$$

It is seen that the data for the  $\text{Mn}^{4+}$ -doped  $\alpha\text{-Al}_2\text{O}_3$  are in very good agreement with the dependence described by Eq. 6. Thus, the values of Racah and crystal-field related parameters for

$\text{Mn}^{4+}$  ions in  $\alpha\text{-Al}_2\text{O}_3$  obtained using the ZPL approach, conform to those of  $\text{Mn}^{4+}$  in other oxide compounds and also to those of  $\text{Cr}^{3+}$  in this host lattice.



**Fig. 5. Dependence of energy of the  $Mn^{4+} 2E_g$  state on the nephelauxetic parameter  $\beta_1$  for oxide compounds. The solid line represents the least-squares fit of the data by expression (6)**

As it was mentioned in Introduction, the present study was also motivated by an interest in obtaining luminescent composites of general formula  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}/\text{GdAG}:\text{Ce}^{3+}$ . The one of advantages of such composite materials is the possibility to tune the emission color by changing weight ratio of the components. Yttrium-gadolinium aluminum garnets ( $\text{Y}_{1-x}\text{Gd}_x$ )<sub>3</sub>Al<sub>5</sub>O<sub>12</sub> (YGdAG) doped with  $\text{Ce}^{3+}$  ions are widely used as the effective yellow phosphors for white LEDs. The restricting factors in using such LEDs for general lighting are a low color rendering index ( $R_a < 80$ ) and a high color temperature ( $> 4000 \text{ K}$ ) that is caused by a deficit of red component in the emission spectrum of  $\text{Ce}^{3+}$  doped YGdAG [28; 29]. GdAG :  $\text{Ce}^{3+}$  is known as a suitable yellow-orange phosphor for warm white LEDs, however, because of the temperature quenching of luminescence, its application is restricted to medium power LEDs [10; 11; 28].

Meanwhile,  $\text{Al}_2\text{O}_3/(\text{Y}_{1-x}\text{Gd}_x)_3\text{Al}_5\text{O}_{12}:\text{Ce}^{3+}$  composite ceramics have been recently proposed as promising materials for high-power LEDs [29]. The positive effect of  $\text{Al}_2\text{O}_3$  addition on thermal stability of the ceramics has been attributed to the high thermal conductivity of  $\alpha\text{-Al}_2\text{O}_3$  ( $32\text{--}35 \text{ W m}^{-1}\text{K}^{-1}$ ) [29]. In the present work, the use of  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}$  as a component of composites had a double purpose: (a) to compensate a deficit of red component in the emission spectrum of  $\text{Gd}_3\text{Al}_5\text{O}_{12}:\text{Ce}^{3+}$  at the

expense of the red emission of  $\text{Mn}^{4+}$  in  $\alpha\text{-Al}_2\text{O}_3$ , (b) to improve the thermal stability of composites owing to the high thermal conductivity of  $\alpha\text{-Al}_2\text{O}_3$ .

As it was mentioned above, the as-prepared  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}$  contains both  $\text{Mn}^{3+}$  (detected by the diffuse-reflectance spectroscopy) and  $\text{Mn}^{4+}$  ions, but only  $\text{Mn}^{4+}$  ions luminesce in  $\alpha\text{-Al}_2\text{O}_3$ . The same conclusion was made previously for Mn-doped  $\alpha\text{-Al}_2\text{O}_3$  obtained by wet chemistry methods [17; 20]. The stabilization of  $\text{Mn}^{4+}$  in  $\text{Al}_2\text{O}_3$  can be provided by the introduction of charge-compensating ions, such as  $\text{Mg}^{2+}$  into the host lattice [8; 19]. In other words, the introduction of  $\text{Mg}^{2+}$  should increase the relative concentration of  $\text{Mn}^{4+}$ . This coincides with the experimental results obtained in the present work. Indeed, from Fig. 6 one can see that the  $\text{Mg}^{2+}$  co-doping into the  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}$  increases the integrated emission intensity by  $\sim 2.9$  times, but it does not result in changes in the  $\text{Mn}^{4+}$  emission spectrum.

The emission spectra of  $\alpha\text{-Al}_2\text{O}_3:\text{Mn}^{4+}, \text{Mg}^{2+}/\text{GdAG}:\text{Ce}^{3+}$  composites with different weight ratios of the components are compared in Fig. 7. One can see that the spectra are superpositions of a broad band extending from 500 to 800 nm and a narrow one at 678 nm.

It is evident that the broad band with a maximum at 585 nm is caused by the  $5d \rightarrow 4f$  transitions of  $\text{Ce}^{3+}$  ions in GdAG [10; 11], whereas the band at 678 nm is caused by the  $\text{Mn}^{4+} 2E_g \rightarrow 4A_{2g}$  transitions in  $\alpha\text{-Al}_2\text{O}_3$ .

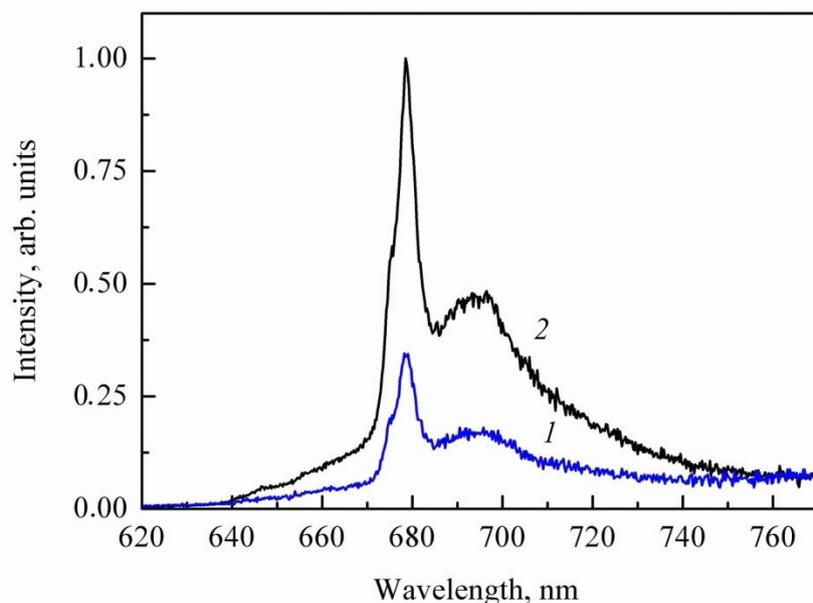


Fig. 6. Effect of the  $\text{Mg}^{2+}$  co-doping into the Mn-doped  $\alpha\text{-Al}_2\text{O}_3$ . The emission spectra of as-prepared  $\alpha\text{-Al}_2\text{O}_3\text{:Mn}^{3+/4+}$  (1) and  $\alpha\text{-Al}_2\text{O}_3\text{:Mn}^{4+},\text{Mg}^{2+}$  (2) were recorded upon excitation at 460 nm

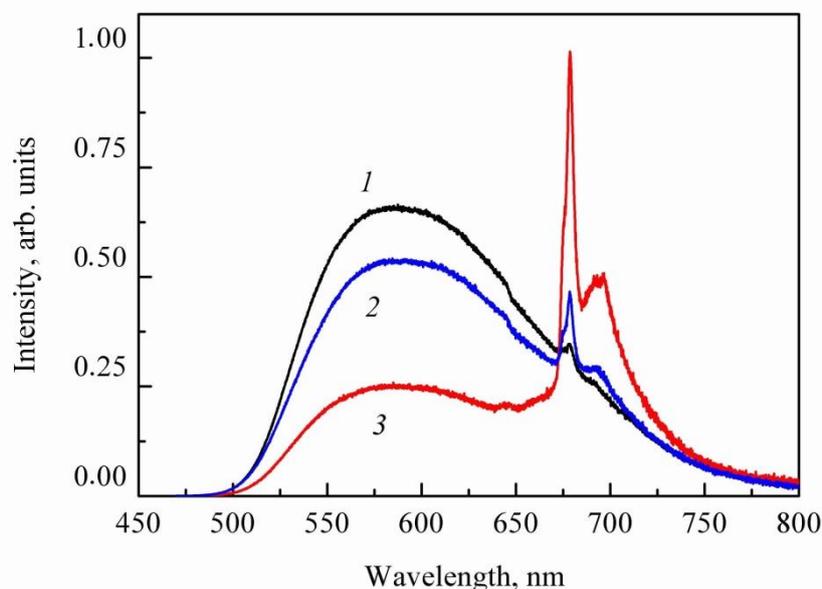


Fig. 7. Comparison of the emission spectra of  $\alpha\text{-Al}_2\text{O}_3\text{:Mn}^{4+},\text{Mg}^{2+}/\text{GdAG:Ce}^{3+}$  composites with different weight ratios (N) of the components: N= 1:2 (1); 1:1 (2); 3:1 (3). The spectra were recorded at 293 K upon excitation with  $\lambda_{\text{exc}} = 460$  nm

It should be noted that the emission color is variable, and it changes from yellow to deep red with increasing content of  $\alpha\text{-Al}_2\text{O}_3\text{:Mn}^{4+},\text{Mg}^{2+}$ . The luminescence quantum efficiencies (QE) of the composites were determined as it was described in Ref. [30] using a commercial  $\text{Y}_3\text{Al}_5\text{O}_{12}\text{:Ce}^{3+}$  phosphor for LEDs with  $\text{QE} = 0.90$  as a standard. For the excitation at 460 nm, depending on the weight ratio of the components the QE value was found to vary in the range from 0.52 to  $\sim 0.60$ . These values are comparable with that ( $\text{QE} = 0.58 \pm 0.04$ ) of the  $\text{Ce}^{3+}$  emission in  $\text{GdAG:Ce}^{3+}$  [11], but they are larger than that ( $\text{QE} = 0.46$ ) reported for  $\text{Mn}^{4+}$  ions in  $\alpha\text{-Al}_2\text{O}_3\text{:Mn}^{4+},\text{Mg}^{2+}$

ceramic phosphors [8]. The luminescence efficiency of the composites can be enhanced by optimization of the synthesis procedure,  $\text{Mn}^{4+}$ -concentration in  $\alpha\text{-Al}_2\text{O}_3$  and composition modification, but even at this stage quite efficient luminescent materials were obtained.

### Conclusions

The  $\text{Mn}^{4+}$ -doped  $\alpha\text{-Al}_2\text{O}_3$  was obtained by combustion method. Based on the results of luminescence measurements, the crystal-field strength ( $Dq$ ) and Racah parameters ( $B$ ,  $C$ ) for  $\text{Mn}^{4+}$  ions in  $\alpha\text{-Al}_2\text{O}_3$  were determined using the ZPL approach. The obtained values of  $Dq$  ( $1898\text{ cm}^{-1}$ ) and nephelauxetic parameter  $\beta_1$

(0.95) for  $Mn^{4+}$  ions in the  $\alpha-Al_2O_3$  are conformed to those of  $Mn^{4+}$  in other oxide compounds and also to those of  $Cr^{3+}$  in the host lattice. In particular, the  $Dq/B$  value of 2.77 indicates a strong crystal field surrounding of the  $Mn^{4+}$  ions in the  $\alpha-Al_2O_3$  lattice. The composites of general formula  $\alpha-Al_2O_3:Mn^{4+}, Mg^{2+}/Gd_3Al_5O_{12}:Ce^{3+}$  were also prepared. These materials are shown to exhibit the intense broadband emission with

maxima at about 585 and 678 nm. The emission color is variable and it changes from yellow to deep red with increasing content of  $\alpha-Al_2O_3:Mn^{4+}, Mg^{2+}$ . Moreover, the luminescence quantum efficiency of the composites was found as high as 0.60. One can expect that the luminescence efficiency of the composites can be enhanced by optimization of the synthesis conditions and  $Mn^{4+}$ -concentration in  $\alpha-Al_2O_3$ .

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